

INTRODUCTION

Centennial scale variability is ubiquitous in historical records of temperature and proxy records in sediments and ice cores. As the slow component of the Earth's climate system, the ocean circulation is a potential candidate for generating oscillations on such long time scales. We developp this idea through the analysis of the stability of the ocean circulation in a hierarchy of simplified ocean models (one- and two-dimensional), using linear stability analysis and density variance budgets in order to better understand the oscillation mechanism.

LATITUDE-DEPTH MODEL

 $\partial_t T = -J(\psi, T) + K_H \partial_u^2 T + K_V \partial_z^2 T + \mathcal{C}$, Temperature evolution $\partial_t S = -J(\psi, S) + K_H \partial_{\psi}^2 S + K_V \partial_z^2 S + \mathcal{C}$, Salinity evolution $\kappa v = -\rho_0^{-1} \partial_y P$, Linear friction equation (Wright and Stocker, 1991) $\partial_z P = -\rho q$, Hydrostatic equation $\partial_{\mu}v + \partial_z w = 0$, Continuity equation Pdz = 0, Baroclinicty condition $\rho = \rho_0 \left[1 - \alpha (T - T_0) + \beta (S - S_0) \right], \text{ Density equation}$ $K_V \partial_z T(y) = \gamma h_m \left(T^*(y) - T(y) \right)$, Temperature boundary condition $K_V \partial_z S(y) = FW(y)$, Salinity boundary condition

Our control parameter for the direct integration experiments is the freshwater flux intensity (F_0) . We explored the range 70–100 cm yr⁻¹. After a Hopf bifurcation (around 79 cm yr⁻¹ without convection and 92 cm yr⁻¹ with convection) the direct integration of our model reveals centennial variability, this variability persists with and without convection (in the following presentation we deal with the case without convection).



Temperature, salinity and overturning averaged over the period; time series of the maximum of the overturning streamfunction. There is perpetual oscillation for the case without convection at $F_0 = 80 \text{ cm yr}^{-1}$ of freshwater flux intensity.

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ON THE MECHANISM OF THE CENTENNIAL THERMOHALINE OSCILLATIONS

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is 171 yr with a growing time scale of 206 yr. growing time scale.



$$X_R \to X_I \to -X_R \to -X_I \to X_R$$
, with $X = \{T', S' \in V \}$
evolution are very similar to the nor
 $\mathbf{HOWARD} \mathbf{NIALKIISIC}$

Malkus, 1972).

$$\partial_t T + \omega \partial_\phi T = r_T \left(T_0 G^T \left(\phi \right) - T \right)$$
$$\partial_t S + \omega \partial_\phi S = -\frac{F_0 S_0}{h} G^S \left(\phi \right)$$
$$\omega = -\int_0^{2\pi} k \left(-\alpha T + \beta S \right) \sin \left(\phi \right) d\phi$$

propagation and surface intensification.



for both).



Variance term	nonlinear	linear	malkus
$- < \overline{\rho' J\left(\psi', \overline{\rho} ight)} > / < \overline{\rho'^2} > 0$	0.000	0.001	-0.001
$<\overline{D'_{ ho} ho'}>/<\overline{ ho'^2}>$	-0.070	-0.072	-0.002
$\alpha^2 \rho_0^2 < \overline{F_T' T'} > / < \overline{\rho'^2} >$	-0.026	-0.029	-0.005
$-\alpha\beta < \overline{F_T'S'} > / < \overline{\rho'^2} >$	0.096	0.094	0.008
$-\alpha\rho_0 < \overline{F_T'\rho'} > / < \overline{\rho'^2} >$	0.070	0.065	0.003
$<\overline{\partial_t \frac{ ho'^2}{2}}>/<\overline{ ho'^2}>$	0.000	-0.006	0.000